

Non-linear far-infrared spectroscopy with an FEL

Thomas Dekorsy

Institute of Ion Beam Physics and Materials Research
Forschungszentrum Rossendorf, Dresden, Germany

The logo for WIRMS 2003 features the text "WIRMS 2003" in a bold, three-dimensional, wood-grained font. The letters are blocky and have a textured surface that resembles wood grain. The text is slightly shadowed, giving it a 3D appearance as if it's floating or standing on a surface.

July 8-11, 2003, Lake Tahoe, CA

Outline

- ◆ Nonlinear optics in the THz frequency range
- ◆ Nonlinear susceptibility of III-V semiconductors
- ◆ Investigation of the nonlinear susceptibility of GaAs below the Reststrahlen band with an FEL
- ◆ Conclusion

Nonlinear optics in the THz frequency range I

- ◆ Nonlinear optics (frequency mixing, harmonic generation,...) well established in the wavelength range $> 20 \mu\text{m}$ to VIS/UV
- ◆ In the THz frequency range (1-10 THz, 300-30 μm) no convenient, high power light sources are available for NLO (“THz gap”)
- ◆ NLO experiments at THz frequencies are based on
 - ↳ frequency mixing of visible with FIR gas laser/microwaves ¹⁾
 - ↳ FIR gas lasers operating at discrete frequencies ²⁾
 - ↳ tunable free-electron lasers
- ◆ Recent developments:
 - ↳ THz quantum cascade lasers ³⁾
 - ↳ DFG/OR with intense fs pulses ⁴⁾
 - ↳ DFG coherent synchrotron THz radiation ⁵⁾

1) Faust and Henry, PRL **17**, 1265 (1966).

2) Mayer, Keilmann, PRB **33**, 6954 (1986), review: Keilmann, Infrared Phys. **31**, 373 (1991).

3) Köhler et al., Nature **417**, 156 (2002).

4) e.g. Reimann et al., Optics Lett. **28**, 471 (2003).

5) Carr et al., Nature **420**, 153 (2002).

Nonlinear optics in the THz frequency range II

- ◆ What is interesting in NLO at THz frequencies?
 - ↳ Many elementary low-energy excitations lie in the THz frequency range, e.g. phonon-polaritons, plasmons, gap energies of superconductors, collective vibrations of biomolecules etc...
 - ↳ Their coupling to radiation gives rise to resonant enhancement of NLO processes
 - ↳ Possibility to obtain information on nonlinear susceptibility at resonances
- ◆ Here: investigation of second harmonic generation below the optical phonon resonance in GaAs¹⁾.
 - ↳ detailed information on nonlinear terms of the lattice potential

1) Dekorsy et al. Phys. Rev. Lett. **90**, 055508 (2003).

Theory: 2nd order nonlinear susceptibility

- ◆ General description of $\chi^{(2)}$ in the vicinity of a phonon-polariton resonance¹⁾:

$$\chi^{(2)}(\omega_3, \omega_2, \omega_1) = \chi_E^{(2)} \left[1 + C_1 \left(\frac{1}{D(\omega_1)} + \frac{1}{D(\omega_2)} + \frac{1}{D(\omega_3)} \right) + C_2 \left(\frac{1}{D(\omega_1)D(\omega_2)} + \frac{1}{D(\omega_1)D(\omega_3)} + \frac{1}{D(\omega_2)D(\omega_3)} \right) \right] \\ \left[+ C_3 \left(\frac{1}{D(\omega_1)D(\omega_2)D(\omega_3)} \right) \right]$$

with
$$D(\omega) = 1 - \left(\frac{\omega}{\omega_{TO}} \right)^2 - \frac{i\omega\gamma}{\omega_{TO}^2}$$

$\chi_E^{(2)}$: pure electronic part

ω_{TO} : TO phonon frequency

γ : phonon damping constant

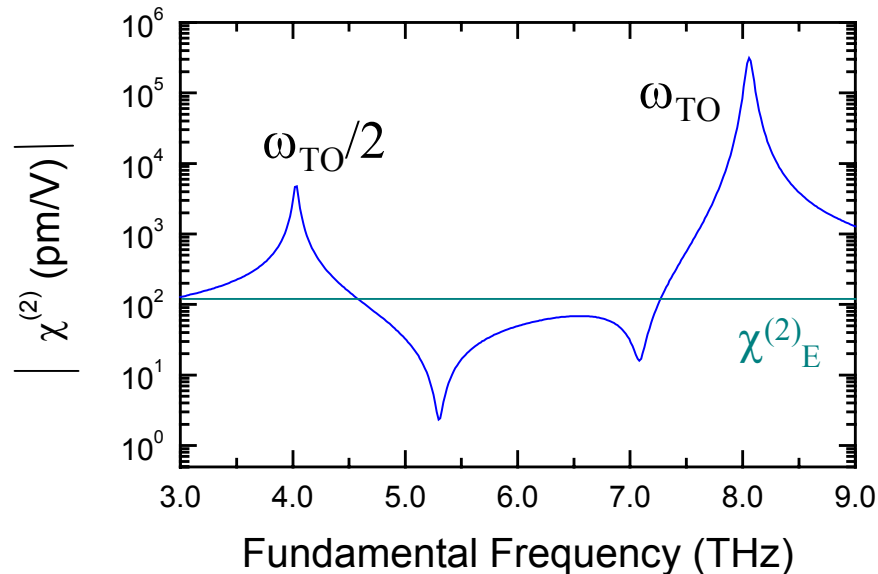
C_1 : Faust-Henry coefficient (Pockels effect, frequency mixing, Raman I_{TO}/I_{LO})

C_2 : second-order lattice dipole moment

C_3 : third-order lattice potential anharmonicity

1) Flytzanis, Phys. Rev. B **6**, 1264 (1972).

$\chi^{(2)}(\omega_3=2\omega_1)$ of GaAs in the THz range



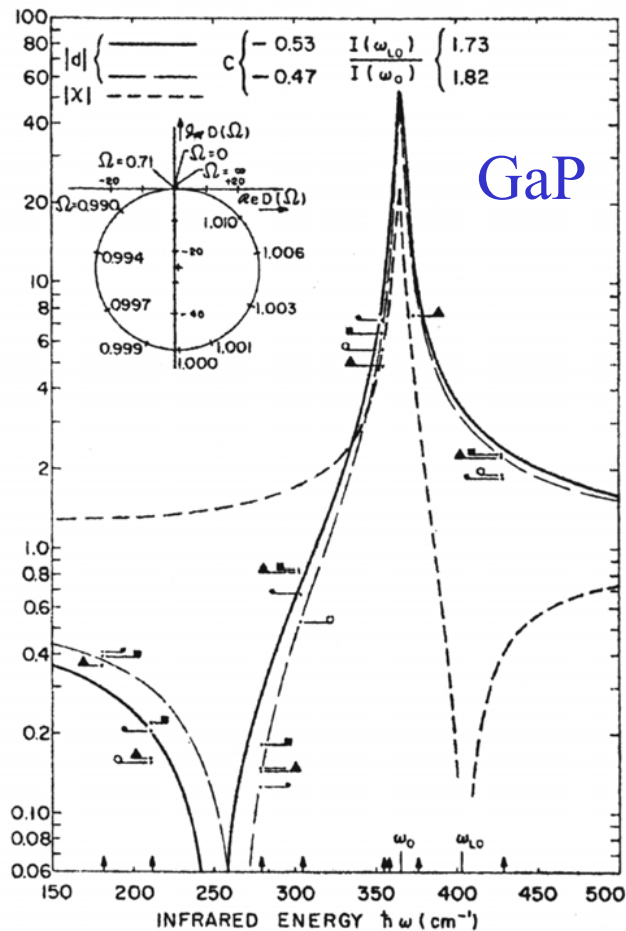
- ◆ $\omega_{\text{TO}} = 8.05$ THz
- ◆ $\chi^{(2)}_{\text{E}} = 132$ pm/V
- ◆ $C_1 = -0.59$ (Faust-Henry coeff.)
- ◆ C_2, C_3 not accurately known:
theory: $C_2 = 0.14$, $C_3 = -0.07$,
 $3C_2 + C_3 = 0.35$;
experiment $3C_2 + C_3 = 0.39$

- ◆ strong resonance at ω_{TO} , exceeding $\chi^{(2)}_{\text{E}}$ by factor of 2100
↳ observed before in GaP¹⁾
- ◆ resonance at $\omega_{\text{TO}}/2$, exceeding $\chi^{(2)}_{\text{E}}$ by a factor of 33
- ◆ 2 minima in $\chi^{(2)}$ between $\omega_{\text{TO}}/2$ and ω_{TO}
↳ resonance and minima not observed before
↳ observation would allow determination of C_i 's without need for determination of absolute intensities

1) Barmentlo et al. Phys. Rev. A **50**, R14 (1994)

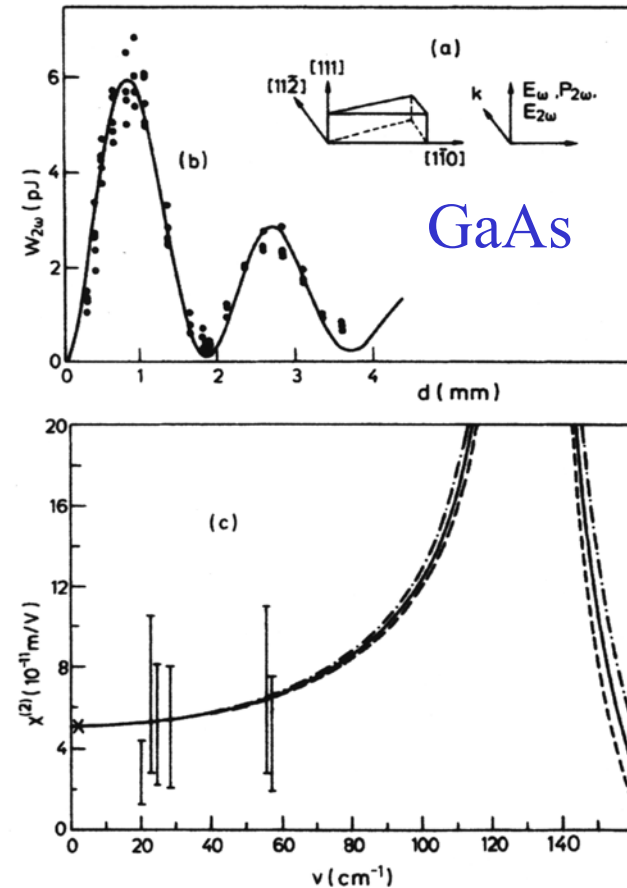
Previous nonlinear optical determinations of C_i 's

frequency mixing of 632 nm + (20-55) μm



Faust et al., PR **173**, 781 (1968)

SHG with CH₃F laser (175-500 μm)

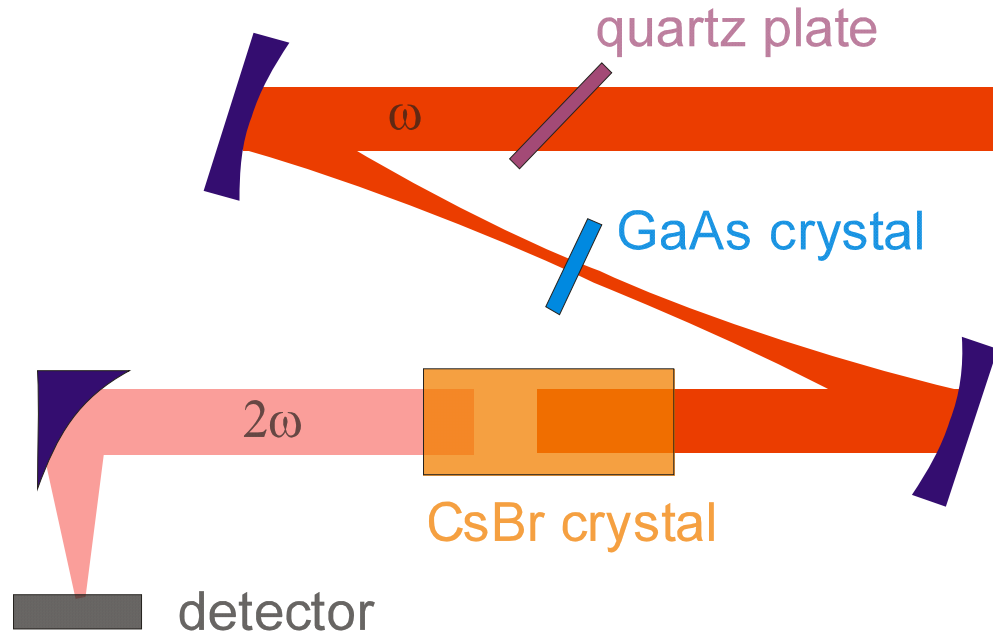


Mayer & Keilmann, PRB **33**, 6954 (1986)

Measurement of $\chi^{(2)}$ with a free-electron laser

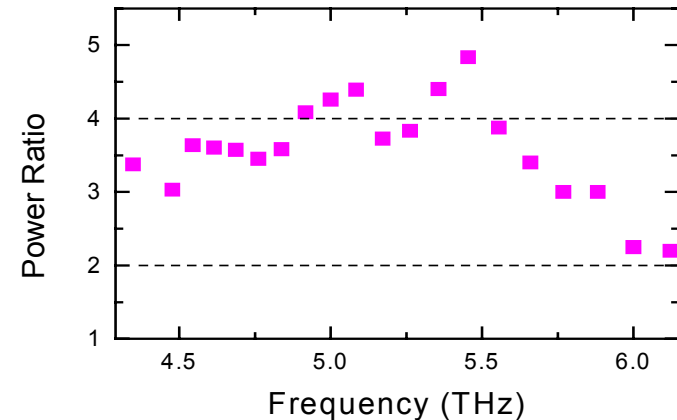
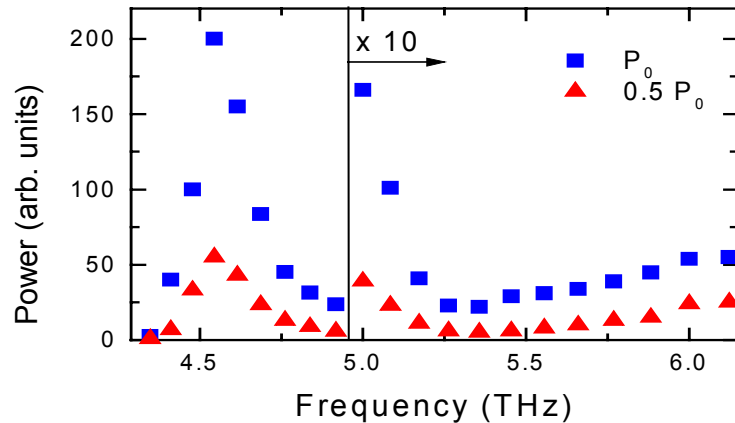
- ◆ free-electron laser FELIX, Rijnhuizen, NL
- ◆ tunable in the THz frequency range (1.2 THz-70 THz)
- ◆ macro-pulses with 10 Hz repetition rate with 100 micro-pulses at 25 MHz
- ◆ picosecond pulse width, 0.2 THz spectral width
- ◆ 4 - 8 μJ energy per pulse, average power 40-80 mW
- ◆ focused on $\sim 500 \mu\text{m}$ spot \Rightarrow 30 kV/cm electric field

Experimental set-up



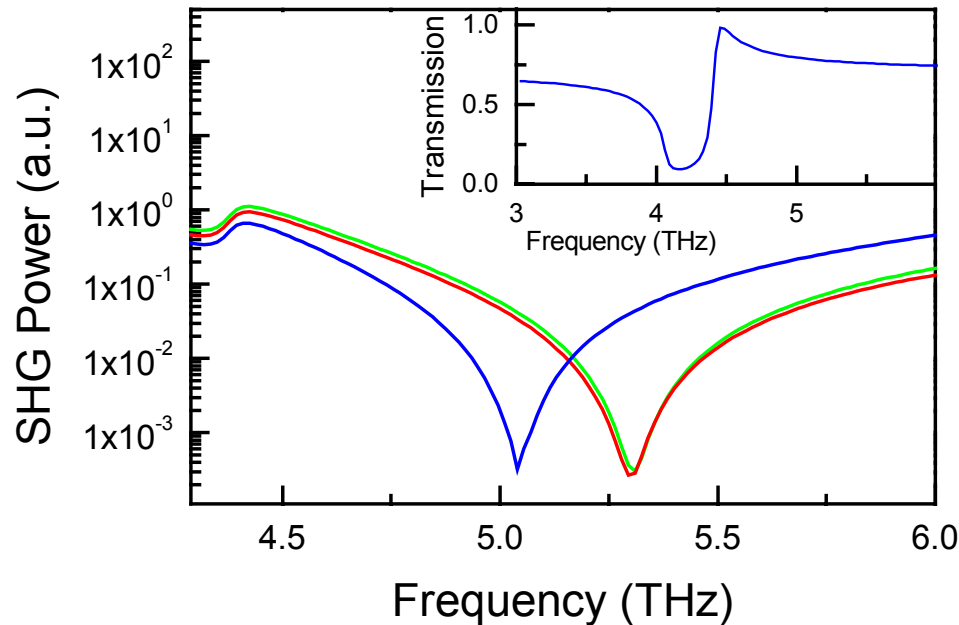
- ◆ 2.7 mm quartz plate: cleans FEL pulses from higher harmonics (fundamental > 4.3 THz) better than 3.5×10^{-7}
- ◆ 2×2.7 cm CsBr crystal: blocks FEL fundamental before detector between 4 THz to 5.5 THz better than 10^{-19} , losses for SHG of 66 %
- ◆ high sensitivity helium cooled Ge:Ga detector
- ◆ whole set-up evacuated

Experimental result on 18 μm thick (211) oriented GaAs crystal



- ◆ Detection of SHG power for two different incident powers
- ◆ Power ratio confirms sensitivity of set-up to SHG between 4.3 THz to 5.6 THz
- ◆ Strong maximum at 4.5 THz ($>\omega_{\text{TO}}/2=4.0$ THz)
- ◆ Minimum at 5.3 THz

Theoretical modeling of observed zero-crossing frequency



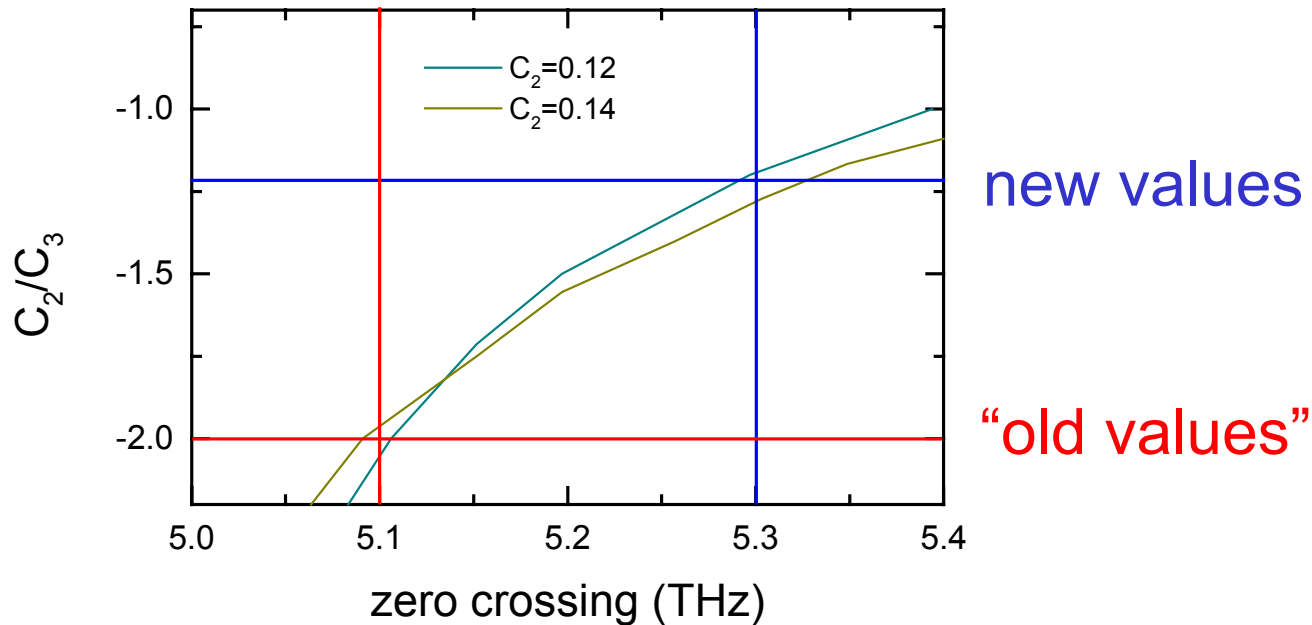
$3C_2+C_3=0.35$, $C_2/C_3= -2.0$
(old theoretical values)

$3C_2+C_3=0.35$, $C_2/C_3= -1.23$

$3C_2+C_3=0.39$, $C_2/C_3= -1.3$
(fits to our experiment)

- ◆ calculation of SHG considering coherence length, absorption, Fresnel coefficients for outcoupling
- ◆ observed zero-crossing point of $\chi^{(2)}$: contribution from phonon interaction through the third-order lattice potential anharmonicity vs second order lattice dipole moment is significantly larger than calculated by theory

Uncertainty of determined C_i 's



- ◆ error in determination of zero-crossing of $\chi^{(2)} < 0.1$ THz
- ◆ ambiguity in determination of C_2 and C_3
- ◆ but: ratio of C_2/C_3 has to be significantly changed

Conclusion

- ◆ Investigation of dispersion of second order nonlinear susceptibility in GaAs with a free-electron laser
- ◆ First observation of resonance close to half the phonon frequency (4.5 THz)
- ◆ Determination of zero-crossing point of $\chi^{(2)}$ at 5.3 THz
- ◆ New values for higher order lattice potentials to the nonlinear susceptibility: contribution from phonon interaction through the third-order lattice potential anharmonicity vs second order lattice dipole moment is significantly larger than calculated by theory¹⁾
- ◆ New method for determination of these parameters in other materials

1) Dekorsy et al. Phys. Rev. Lett. **90**, 055508 (2003).

Acknowledgments

Vladimir Yakovlev	Institute of Spectroscopy, Russian Academy of Science, Troitsk, Russia
Wolfgang Seidel	Forschungszentrum Rossendorf, Dresden, Germany
Manfred Helm	Forschungszentrum Rossendorf, Dresden, Germany
Fritz Keilmann	Max-Planck-Institut für Biochemie, Martinsried, Germany

Team of FELIX FEL facility for their support.

Work supported under the “Access to research infrastructure” action of the European Community